

**Effective dynamics for low-amplitude transient elastic waves
in a 1D periodic array of non-linear interfaces**

Cédric Bellis^{1,*}, Bruno Lombard¹, Marie Touboul¹, Raphaël Assier^{2,2}

¹Aix Marseille Univ, CNRS, Centrale Marseille, LMA, Marseille, France

²Department of Mathematics, The University of Manchester, Oxford Road, Manchester, UK

*Email: bellis@lma.cnrs-mrs.fr

Abstract

This presentation focuses on the time-domain propagation of elastic waves through a 1D periodic medium that contains non-linear imperfect interfaces, i.e. interfaces exhibiting a discontinuity in displacement and stress governed by a non-linear constitutive relation. In this context, we investigate transient waves with both low-amplitude and long-wavelength, and aim at deriving homogenized models that describe their effective motion.

Keywords: Homogenization – Correctors – Imperfect interfaces – Non-linear waves – Time-domain numerical simulations

1 Objectives

The array of interfaces considered is generated by a, possibly heterogeneous, cell repeated periodically and bonded by interfaces that are associated with transmission conditions of non-linear “spring-mass” type. More precisely, the imperfect interfaces are characterized by a linear dynamics but a non-linear elasticity law. The latter is not specified at first and only key theoretical assumptions are required. To establish an effective model, the two-scale asymptotic homogenization method is deployed, up to the first-order. To begin, an effective model is obtained for the leading zeroth-order contribution to the microstructured wavefield. It amounts to a wave equation with a non-linear constitutive stress-strain relation that is inherited from the behavior of the imperfect interfaces at the microscale. The next first-order corrector term is then shown to be expressed in terms of a cell function and the solution of a linear elastic wave equation. Without further hypothesis, the constitutive relation and the source term of the latter depend non-linearly on the zeroth-order field, as does the cell function. Combining these zeroth- and first-order models leads to an approximation of both the macroscopic behavior of the microstructured wavefield and its small-

scale fluctuations within the periodic array.

2 Setting: microstructured configuration

We consider the propagation of transient waves in a 1D periodic elastic medium containing imperfect interfaces. The latter have spacing h and, for simplicity but with no loss of generality, we consider that they are located at $X_n = nh$ with $n \in \mathbb{Z}$. The elastic medium is supposed to be h -periodic and linear elastic with mass density $\rho_h(X)$ and Young’s modulus $E_h(X)$. Given a source term F , the displacement field U_h is governed by the time-domain wave equation

$$\rho_h(X) \frac{\partial^2 U_h}{\partial t^2}(X, t) = \frac{\partial \Sigma_h}{\partial X}(X, t) + F(X, t) \quad (1)$$

where

$$\Sigma_h(X, t) = E_h(X) \frac{\partial U_h}{\partial X}(X, t),$$

with Σ_h being the stress field. Moreover, the interfaces are assumed to be characterized by the interface *mass* and *rigidity* parameters M and K , respectively, together with the, possibly *non-linear*, constitutive relation \mathcal{R} , so that the following transmission conditions apply at any interface point X_n , see [1–4]:

$$\begin{cases} M \left\langle \left\langle \frac{\partial^2 U_h}{\partial t^2}(\cdot, t) \right\rangle \right\rangle_{X_n} = \llbracket \Sigma_h(\cdot, t) \rrbracket_{X_n} & (2a) \\ \llbracket \Sigma_h(\cdot, t) \rrbracket_{X_n} = K \mathcal{R}(\llbracket U_h(\cdot, t) \rrbracket_{X_n}), & (2b) \end{cases}$$

where, for any function $g(X)$, we define the jump and mean operators $\llbracket \cdot \rrbracket_{X_n}$ and $\langle \langle \cdot \rangle \rangle_{X_n}$ as

$$\begin{aligned} \llbracket g \rrbracket_{X_n} &= g(X_n^+) - g(X_n^-), \\ \langle \langle g \rangle \rangle_{X_n} &= \frac{1}{2}(g(X_n^+) + g(X_n^-)). \end{aligned} \quad (3)$$

In addition, both the displacement U_h and the stress field Σ_h are continuous on the open intervals (X_n, X_{n+1}) .

3 Main homogenization results

We now consider a reference wavelength λ^* and introduce the following parameters

$$k^* = \frac{2\pi}{\lambda^*} \quad \text{and} \quad \eta = hk^*, \quad (4)$$

k^* being the reference wavenumber. In this study it is assumed that $\eta \ll 1$ and that the source term F is of relatively *low-amplitude* (an issue that will be discussed). The objective is to derive an effective dynamical model, up to the first-order, for the waves propagating in the periodic interface array considered. More precisely, we seek an approximation $U^{(1)}$ of the solution U_h to (1-2) of the form:

$$U_h(X, t) = U^{(1)}(X, t) + o(h).$$

The main results of this study is that the sought-after approximation is given by

$$U^{(1)}(X, t) = U_0(X, t) + hU_1(X, t), \quad (5)$$

where the zeroth-order field U_0 in (5) is continuous and is solution of the problem

$$\rho_{\text{eff}} \frac{\partial^2 U_0}{\partial t^2}(X, t) = \frac{\partial \Sigma_0}{\partial X}(X, t) + F(X, t)$$

with

$$\Sigma_0(X, t) = \mathcal{G}_{\text{eff}}(\mathcal{E}_0(X, t)).$$

Here, $\mathcal{E}_0 = \partial U_0 / \partial X$ and \mathcal{G}_{eff} is an effective strain-stress relation that is local and, generally speaking, non-linear, while ρ_{eff} is an effective mass density. Moreover, the first-order corrector field U_1 in (5) can be written as

$$U_1(X, t) = \bar{U}_1(X, t) + \mathcal{P}(y, \mathcal{E}_0(X, t)) \mathcal{E}_0(X, t)$$

with $y = (X - nh)/h$ for $X \in (nh, (n+1)h)$ and where the *cell function* \mathcal{P} is, generally speaking, a non-linear function of \mathcal{E}_0 . The *mean field* \bar{U}_1 is solution to the *linear* problem:

$$\rho_{\text{eff}} \frac{\partial^2 \bar{U}_1}{\partial t^2}(X, t) = \frac{\partial \bar{\Sigma}_1}{\partial X}(X, t) + \mathcal{S}(U_0(X, t))$$

with

$$\bar{\Sigma}_1(X, t) = \mathcal{G}'_{\text{eff}}(\mathcal{E}_0(X, t)) \bar{\mathcal{E}}_1(X, t),$$

where $\bar{\mathcal{E}}_1 = \partial \bar{U}_1 / \partial X$, while both the *parameter* $\mathcal{G}'_{\text{eff}}(\mathcal{E}_0(X, t))$, which is the derivative of \mathcal{G}_{eff} , and the source term $\mathcal{S}(U_0(X, t))$ depend explicitly on the zeroth-order field, locally in space

and time, and in a non-linear fashion.

Particularizing for a prototypical non-linear interface law and in the cases of a homogeneous periodic cell and a bilaminated one, the behavior of the obtained models will then be illustrated on a set of numerical examples and compared with full-field simulations. Both the influence of the dominant wavelength and of the wavefield amplitude will be investigated numerically, as well as the characteristic features related to non-linear phenomena.

References

- [1] J. D. Achenbach and A. N. Norris. Loss of specular reflection due to nonlinear crack-face interaction. *Journal of Nondestructive Evaluation*, 3(4):229–239, 1982.
- [2] S.C. Bandis, A.C. Lumsden, and N.R. Barton. Fundamentals of rock joint deformation. *International Journal of Rock Mechanics and Mining Sciences & Geomechanics Abstracts*, 20(6):249 – 268, 1983.
- [3] D. Broda, W.J. Staszewski, A. Martowicz, T. Uhl, and V.V. Silberschmidt. Modelling of nonlinear crack-wave interactions for damage detection based on ultrasound—a review. *Journal of Sound and Vibration*, 333(4):1097 – 1118, 2014.
- [4] I. Sevostianov, R. Rodriguez-Ramos, R. Guinovart-Diaz, J. Bravo-Castillero, and F.J. Sabina. Connections between different models describing imperfect interfaces in periodic fiber-reinforced composites. *International Journal of Solids and Structures*, 49(13):1518 – 1525, 2012.